

# Lecture 20

(1)

## Horizontal vorticity equations

x-comp vorticity is  $\frac{\partial w}{\partial y} - \frac{\partial v}{\partial z}$  (call it  $\omega_x$ )

y-comp vorticity is  $\frac{\partial v}{\partial z} - \frac{\partial w}{\partial x}$  (call it  $\omega_y$ )

To get the x-comp vorticity  $eq^n$  ( $eq^n$  for  $\partial \omega_x / \partial \tau$ , take  $\partial / \partial y$  ( $z$   $eq^n$  of motion) and subtract from it  $\partial / \partial z$  ( $y$   $eq^n$  of motion).

To get the y-comp vorticity  $eq^n$  ( $eq^n$  for  $\partial \omega_y / \partial \tau$ , take  $\partial / \partial z$  ( $x$   $eq^n$  of motion) and subtract from it  $\partial / \partial x$  ( $z$   $eq^n$  of motion).

Get:

$$\frac{\partial \omega_x}{\partial \tau} = \text{advection, stretching, Tilting} + \hat{i} \cdot \left( \frac{1}{\rho^2} \nabla \rho \times \nabla p \right)$$

it's large where  $\frac{\partial \rho}{\partial y}$  is large

x-comp baroclinic vorticity generation.

$$\frac{\partial \omega_y}{\partial \tau} = \text{advection, stretching, Tilting} + \hat{j} \cdot \left( \frac{1}{\rho^2} \nabla \rho \times \nabla p \right)$$

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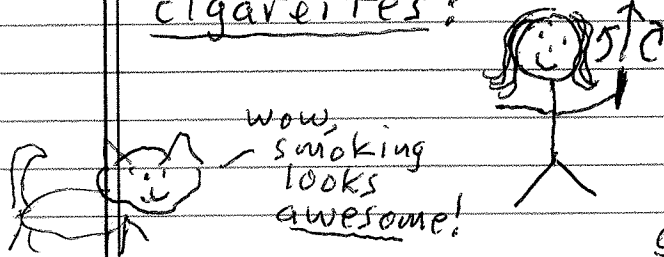
y-comp baroclinic vorticity generation

Flows where the horizontal (x or y) components of baroclinic vorticity generation are important are found across almost all scales in the atmosphere, from microscale to synoptic-scale, for example:

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- Convective clouds, squall lines, supercells, etc.
- leading edge of fronts and  $\tau$ -storm outflows
- gravity waves
- Katabatic flows (flows down cooled slopes)
- anabatic flows ("up heated"),
- sea-breezes
- man-made convective flows arising from:

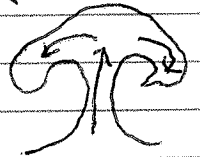
cigarettes:



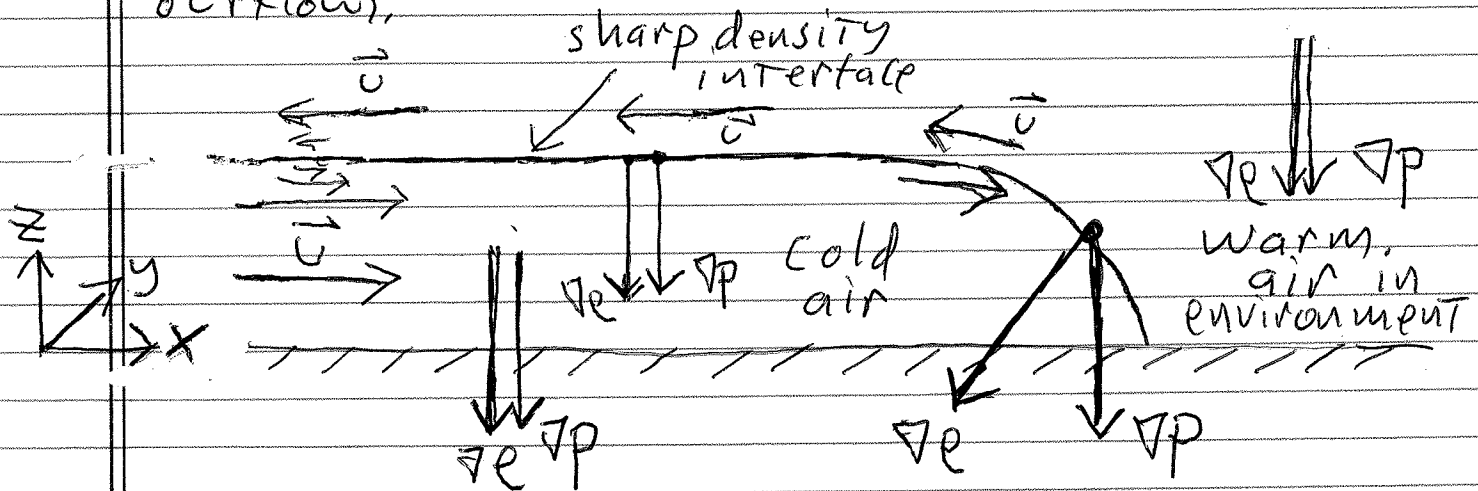
bar-b-ques:



atom bomb blasts:  
(mushroom cloud)



→ Let's focus on one example: a thunderstorm outflow moving in the  $+\hat{i}$  direction (westerly outflow).



What's going on here?

We idealize the outflow as a long blob of cold air displacing relatively warm air in

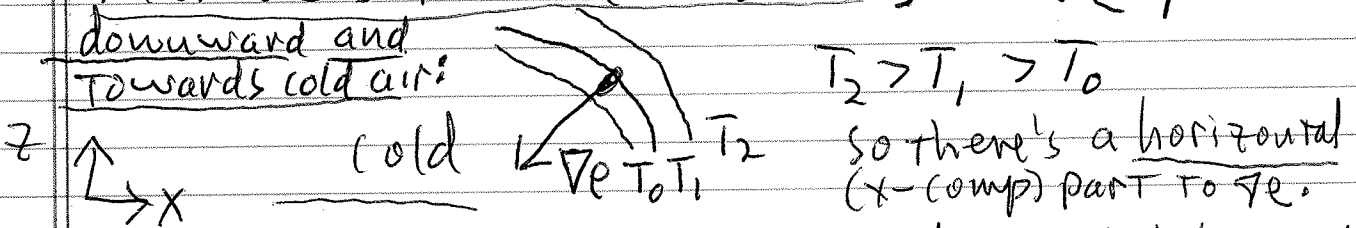
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environment. In environment,  $p$  and  $e$  both increase with decreasing  $z$ , so  $\nabla p$  and  $\nabla e$  both point downward.

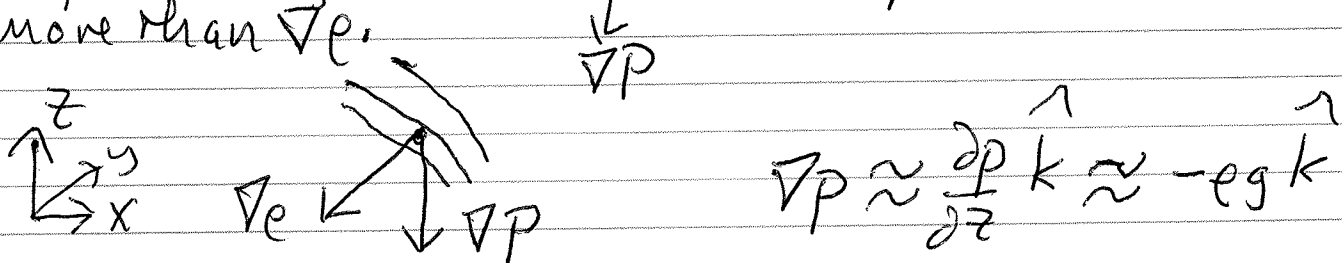
Similarly,  $\nabla p$  and  $\nabla e$  both point downward inside cold pool (except at leading edge).

But something interesting is happening along the leading edge, where there are horizontal changes in  $e$ . Along the

leading edge,  $e$  increases downward and also increases towards cold air, so  $\nabla e$  points downward and



In contrast, the pressure field is still largely hydrostatic, and thus tends to point downward more than  $\nabla e$ .

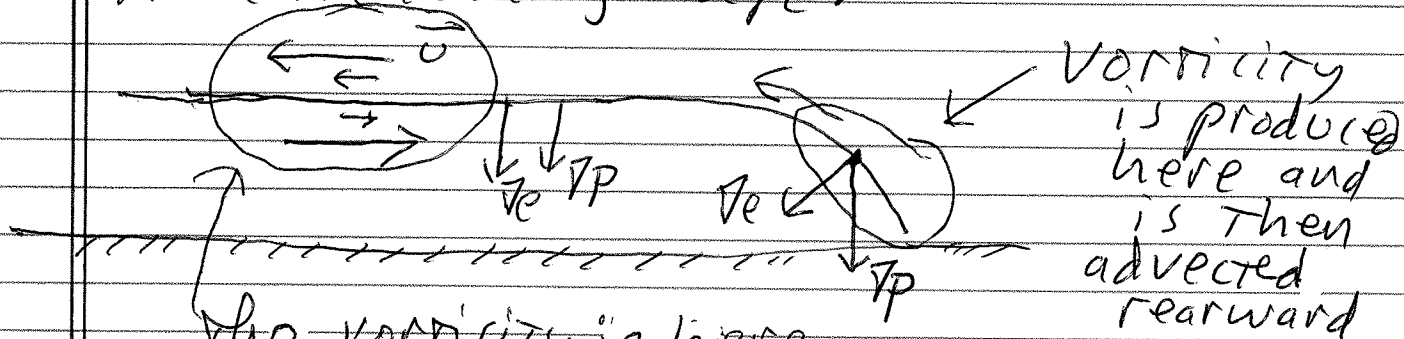


So, on leading edge of interface,  $\nabla e \times \nabla p$  points out-of-page, which is in the minus  $\hat{j}$  direction!

$$\therefore \hat{j} \cdot \left( \frac{1}{\rho^2} \nabla e \times \nabla p \right) = \hat{j} \cdot (-\hat{j} \text{ stuff}) < 0$$

$\therefore \frac{\partial \omega_y}{\partial t} < 0$   $\therefore$  generate negative (out of page) y-comp vorticity.

Note that out-of-page vorticity is what we see along the interface further back from the leading edge:



The vorticity in here was produced baroclinically at the leading edge.

But there is no baroclinic generation going on here because  $\nabla_e$  and  $\nabla_p$  are both pointing nearly straight downward.

### A theory for tornadogenesis (for some tornadoes)

- (i) Horiz vorticity is generated baroclinically at leading edge of  $\tau$ -storm outflow. (baroclinic term in  $x$  or  $y$  comp vorticity equation).
- (ii) This horiz vorticity gets tilted into the vertical by an updraft or a downdraft (+tilting term in  $z$ -comp vorticity eqn). So now we have  $\zeta$ .
- (iii)  $\zeta$  gets stretched (amplified) in a convergent updraft  $\rightarrow$  get a tornado (from stretching term in  $z$ -comp vorticity eqn).

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Scale analysis of vertical vorticity eq<sup>n</sup> for mid-latitude synoptic-scale flows.

Characteristic scales of motion for synoptic-scale flows (based on typical observations):

Horizontal length scale  $L \sim 1000 \text{ km} = 10^6 \text{ m}$

Vertical length scale  $H \sim 10 \text{ km} = 10^4 \text{ m}$

Horizontal velocity scale  $U \sim 10 \text{ m/s}$

Vertical velocity scale  $W \sim 0.01 \text{ m/s}$   
[1 cm/s]

Time scale based on horizontal advection

(Time it takes a parcel moving horizontally at speed  $U$  to traverse the distance  $L$ ),  $T_{\text{horiz}} = \frac{L}{U} = 10^5 \text{ s}$   
(about 1 day)

Time scale based on vertical advection

(Time it takes a parcel moving vertically at speed  $W$  to traverse the vertical distance  $H$ ),  $T_{\text{vert}} = \frac{H}{W} = 10^6 \text{ s}$  (about 10 days)

Typical density  
 $\rho \sim 1 \text{ kg/m}^3$

Typical horizontal change in density (across horizontal distance  $L$ ),

$\Delta \rho_{\text{horiz}} \sim 0.01 \text{ kg/m}^3$  [about 1% of  $\rho$ ]

Typical horizontal change in pressure (across horiz distance  $L$ ),

$\Delta P_{\text{horiz}} \sim 10 \text{ mb} = 1000 \text{ Pa}$  [1 mb = 100 Pa]