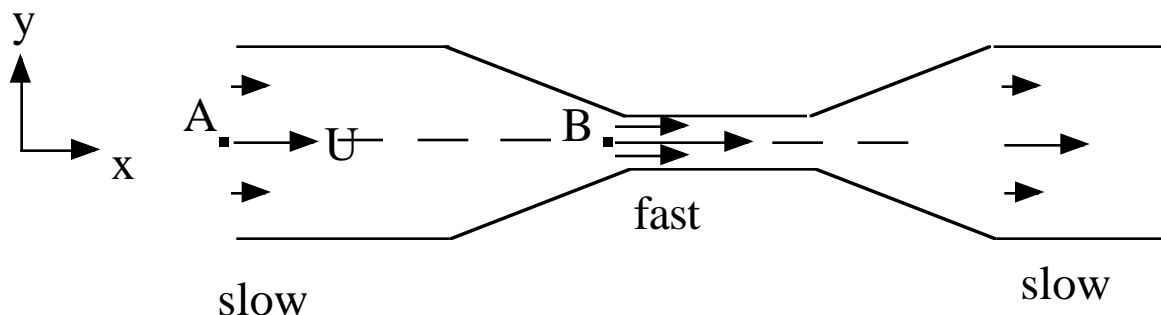


METR 5113, Advanced Atmospheric Dynamics I  
 Alan Shapiro, Instructor  
 Friday, 12 Rocktober 2018 (lecture 22)

**1 handout: answers to problem set 3.**

**Applications of Bernoulli's eqns**

Example 1: Constriction in a channel (flow through a canyon)



Assume steady state and  $\rho = \text{const.}$  Use steady state version of B eqn (where vort may or may not be zero):

$$\frac{q^2}{2} + \frac{p}{\rho} + gz = C \quad (\text{along a streamline})$$

Consider flow along central streamline (through A and B). Far upstream of constriction (at A)  $\vec{u} = U\hat{i}$ , so  $q_A = U$ .

$$\therefore \frac{U^2}{2} + \frac{p_A}{\rho} + gz_A = C$$

At the constriction (B):

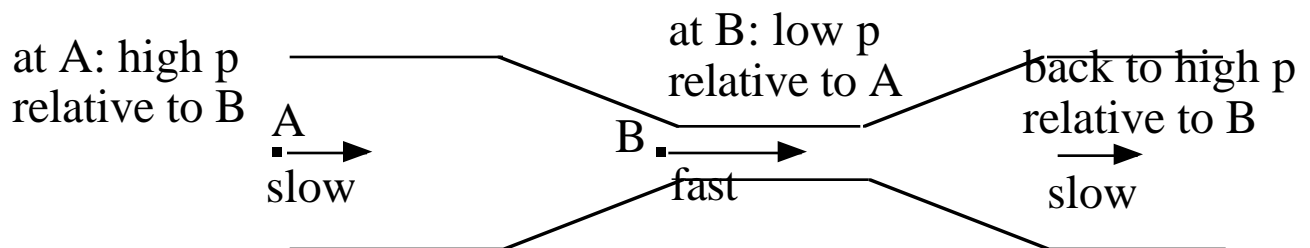
$$\frac{q_B^2}{2} + \frac{p_B}{\rho} + gz_B = C$$

Same C in both eqns because we're on same streamline.

$$\therefore \frac{U^2}{2} + \frac{p_A}{\rho} + \boxed{gz_A} = \frac{q_B^2}{2} + \frac{p_B}{\rho} + \boxed{gz_B} \quad \begin{array}{l} \text{cancellation} \\ \text{since } z_A = z_B \end{array}$$

$$\therefore p_A - p_B = \frac{\rho}{2}(q_B^2 - U^2) > 0 \quad \begin{array}{l} \text{fast} \\ \text{slow} \end{array} \quad \text{since } q_B > U \text{ (from mass cons)}$$

$$\therefore p_A > p_B$$

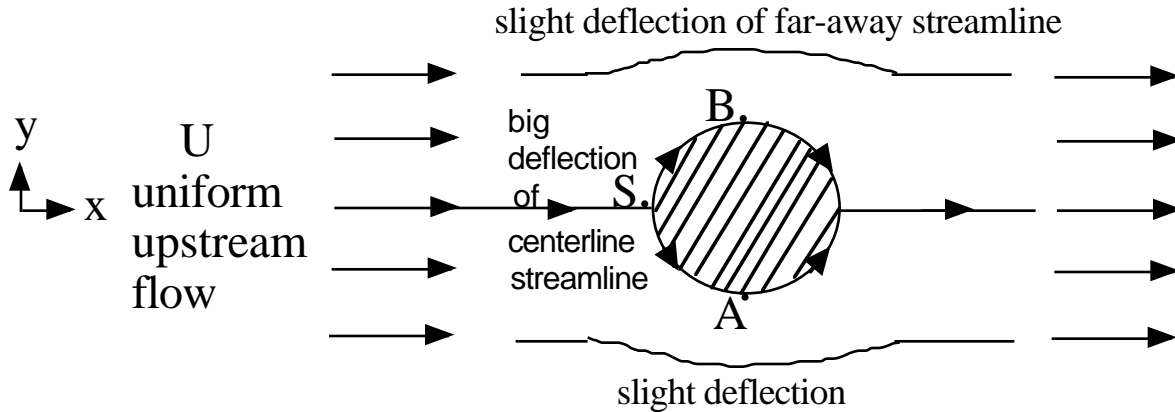


Flow is faster at B than at A (from mass cons eq<sup>n</sup>). So a parcel accelerates from A to B. Bernoulli's eq<sup>n</sup> says this acceleration is associated with a pressure drop.

Get qualitatively same result if you assume uniform flow upstream and use irrotational version of Bernoulli's eqn.

Theoretically, very fast speeds can be associated w/ negative pressures (tensions rather than compressions). In liquids get cavitation (boiling) when vapor pressure is reached. Collapsing vapor bubbles can damage dam spillways, submarine propellers and pumps/pipes.

Example 2 - Steady, 2D, deflection of a uniform flow around a cylinder (e.g. telephone pole). (t-storm updrafts sometimes behave as if they're solid cylinders)



$S$  is a stagnation point -- all velocity comps are 0 at  $S$ . Normal comp at  $S$  is 0 due to impermeability cond<sup>n</sup>. Tangential comp at  $S$  is 0 from symmetry (not from no-slip cond<sup>n</sup>).  $\therefore q_s = 0$ .

From mass cons and streamline deflection (confluence) north and south of cylinder, expect fast flow north and south of cylinder.

Use Bernoulli eq<sup>n</sup> for irrotational flow but w/ unsteady term = 0 since this is a steady state. Assume  $\rho = \text{const}$ .

$$\frac{q^2}{2} + \frac{p}{\rho} + gz = C \quad (\text{same const everywhere})$$

Far upstream (anywhere upstream):

$$\frac{U^2}{2} + \frac{p_{\text{up}}}{\rho} + gz = C$$

At stagnation point at  $S$ :

$$0 + \frac{p_s}{\rho} + gz = C \quad (\text{same const as above})$$

$$\therefore \frac{U^2}{2} + \frac{p_{\text{up}}}{\rho} + \boxed{gz} = \frac{p_s}{\rho} + \boxed{gz} \quad \text{cancellation}$$

$$\therefore p_s = p_{\text{up}} + \rho \frac{U^2}{2}$$

$$\therefore p_s > p_{\text{up}} \quad \therefore \underline{\text{high pressure at stagnation point.}}$$

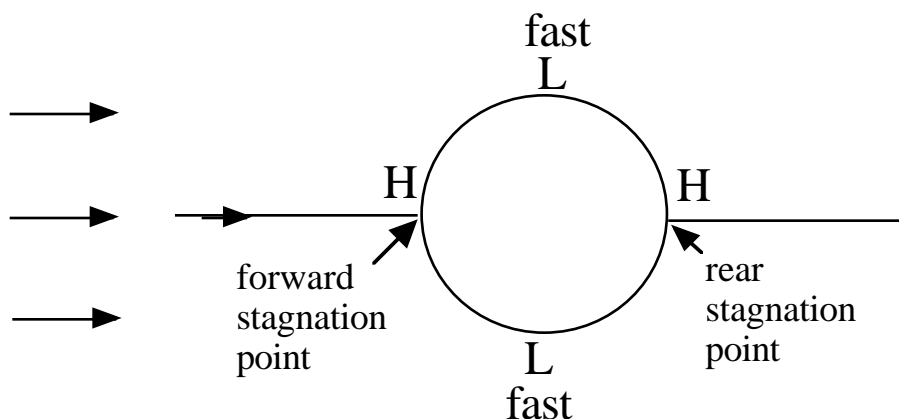
Now look on north point of cylinder:

$$\frac{q_B^2}{2} + \frac{p_B}{\rho} + gz = C$$

$$\therefore \frac{q_B^2}{2} + \frac{p_B}{\rho} + gz = \frac{U^2}{2} + \frac{p_{\text{up}}}{\rho} + gz$$

$$\therefore p_B = p_{\text{up}} + \frac{\rho}{2} \boxed{(U^2 - q_B^2)} \rightarrow \text{negative since } U < q_B$$

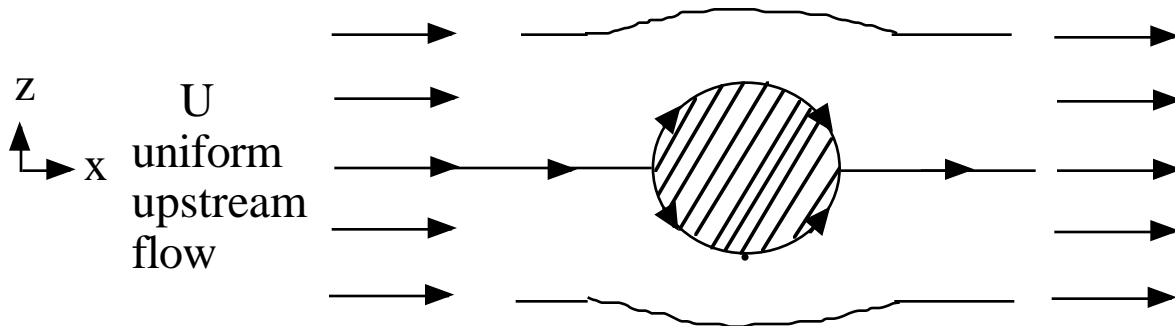
$\therefore p_B < p_{\text{up}} \quad \therefore \underline{\text{low pressure at north point of cylinder}}$  (and at south point)



In real life  $p$  at rear stag point is not as big as  $p$  at forward stag point (due to friction, boundary layer separation in lee of

cylinder, and possible turbulence). End up with a net pressure force on cylinder in direction of flow. However, analysis of forward part of cylinder and upstream still valid (a la Prandtl).

Example 3 - Same as example 2 but consider flow over horizontally-oriented cylinder (or tree branch).



$$\frac{q^2}{2} + \frac{p}{\rho} + gz = C \quad (\text{same const everywhere})$$

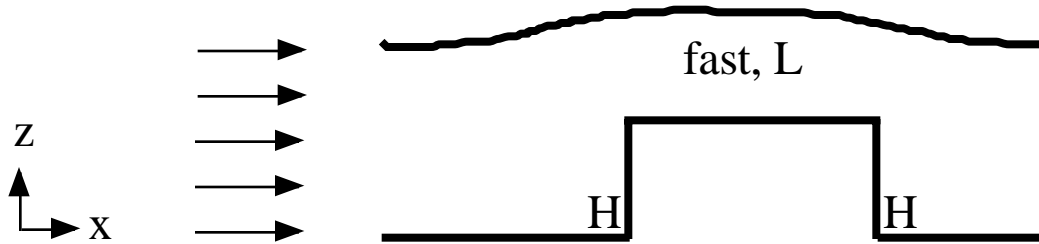
Rewrite using perturbation pressure  $p'$ . [ $p = p' + \bar{p}(z)$   
 $= p' + (\text{const} - \rho gz) = p' - \rho gz + \text{const}$ ].

$$\therefore \frac{p}{\rho} + gz = \frac{p'}{\rho} + \text{const}$$

$$\therefore \frac{q^2}{2} + \frac{p'}{\rho} = C \quad (\text{this } C \text{ differs from old } C)$$

Proceed as in example 2, but work w/ pert pressure  $p'$  instead of full  $p$ . Get low perturbation pressure  $p'$  at top and bottom of cylinder, and high perturbation pressure at stag point.

Example 4: Wind-storm over a Walmart.



Similar to previous example. Can use irrot version of B eqn w/unsteady term set to 0.

$$\frac{q^2}{2} + \frac{p}{\rho} + gz = C \text{ (same const everywhere)}$$

$$\therefore \frac{q^2}{2} + \frac{p'}{\rho} = C$$

Get high  $p'$  at 2 stag points. Get fast flow above roof. There  $q^2$  is really big so  $p'$  is large negative (low  $p$ ) but  $p$  inside building is nearly the same as before storm (it was hydrostatic so  $p'=0$  there).  $\therefore$  net pert pressure force on roof is large and upward -- roof lifts off. Then pressures equalize and roof falls into store.