

METR 5113, Advanced Atmospheric Dynamics I
 Alan Shapiro, Instructor
 Friday, 2 November 2018 (lecture 31)

Ertel's Potential Vorticity Theorem (continued)

The vorticity eqn for an inviscid compressible and baroclinic flow can be written as

$$(B) \quad \frac{D}{Dt} \left(\frac{\vec{\omega}}{\rho} \right) = \left(\frac{\vec{\omega}}{\rho} \cdot \nabla \right) \vec{u} + \frac{1}{\rho^3} \nabla \rho \times \nabla p$$

Now suppose there is some scalar quantity λ that is conserved:

$$\frac{D}{Dt} \lambda = 0 \quad \rightarrow \quad \lambda = \text{const (for a parcel)}$$

λ is a "passive scalar" or "marker" for the flow, e.g., can choose:

$\lambda = \theta$ (pot temp) if flow is isentropic,

$\lambda =$ ozone concentration (for short time scales),

$\lambda =$ concentration of a non-reactive pollutant.

Even if λ is conserved, does not mean that $\nabla \lambda$ is conserved!
 [--> | | | | <-- becomes ||||. Gradient can also rotate, e.g., in presence of a shear]

Look at $\frac{D}{Dt} \nabla \lambda$ (well, one component of it):

$$\frac{D}{Dt} \frac{\partial \lambda}{\partial x_i} = \frac{\partial}{\partial t} \left(\frac{\partial \lambda}{\partial x_i} \right) + u_j \frac{\partial}{\partial x_j} \left(\frac{\partial \lambda}{\partial x_i} \right)$$

$$\begin{aligned}
&= \frac{\partial}{\partial x_i} \left(\frac{\partial \lambda}{\partial t} \right) + u_j \frac{\partial}{\partial x_i} \left(\frac{\partial \lambda}{\partial x_j} \right) \\
&= \frac{\partial}{\partial x_i} \left(\frac{\partial \lambda}{\partial t} \right) + \frac{\partial}{\partial x_i} \left(u_j \frac{\partial \lambda}{\partial x_j} \right) - \frac{\partial u_j}{\partial x_i} \frac{\partial \lambda}{\partial x_j} \\
&= \frac{\partial}{\partial x_i} \left(\frac{\partial \lambda}{\partial t} + u_j \frac{\partial \lambda}{\partial x_j} \right) - \frac{\partial u_j}{\partial x_i} \frac{\partial \lambda}{\partial x_j} \\
&= \frac{\partial}{\partial x_i} \left(\frac{D\lambda}{Dt} \right) - \frac{\partial u_j}{\partial x_i} \frac{\partial \lambda}{\partial x_j} \quad \text{but } \frac{D\lambda}{Dt} = 0 \text{ since } \lambda \text{ is conserved}
\end{aligned}$$

$$\therefore \frac{D}{Dt} \frac{\partial \lambda}{\partial x_i} = - \frac{\partial u_j}{\partial x_i} \frac{\partial \lambda}{\partial x_j} \quad \text{contract with } \omega_i:$$

$$\omega_i \frac{D}{Dt} \frac{\partial \lambda}{\partial x_i} = - \omega_i \frac{\partial u_j}{\partial x_i} \frac{\partial \lambda}{\partial x_j}$$

$$\therefore \bar{\omega} \cdot \frac{D}{Dt} \nabla \lambda = - \nabla \lambda \cdot (\bar{\omega} \cdot \nabla) \bar{u} \quad \div \text{ by } \rho:$$

$$(C) \quad \boxed{\frac{\bar{\omega}}{\rho} \cdot \frac{D}{Dt} \nabla \lambda = - \nabla \lambda \cdot \left(\frac{\bar{\omega}}{\rho} \cdot \nabla \right) \bar{u}}$$

Take $\nabla \lambda \cdot (B)$

$$\nabla \lambda \cdot \frac{D}{Dt} \left(\frac{\bar{\omega}}{\rho} \right) = \nabla \lambda \cdot \left(\frac{\bar{\omega}}{\rho} \cdot \nabla \right) \bar{u} + \frac{1}{\rho^3} \nabla \lambda \cdot (\nabla \rho \times \nabla p)$$

Add this eqⁿ to (C), get:

$$\frac{\bar{\omega}}{\rho} \cdot \frac{D}{Dt} \nabla \lambda + \nabla \lambda \cdot \frac{D}{Dt} \left(\frac{\bar{\omega}}{\rho} \right) = \frac{1}{\rho^3} \nabla \lambda \cdot (\nabla \rho \times \nabla p)$$

$$\frac{D}{Dt} \left(\frac{\bar{\omega}}{\rho} \cdot \nabla \lambda \right) = \frac{1}{\rho^3} \boxed{\nabla \lambda \cdot (\nabla \rho \times \nabla p)}$$

scalar triple product

rhs is 0 if:

(i) density is constant

or (ii) flow is barotropic

or (iii) $\lambda = \lambda(p)$ or $\lambda = \lambda(\rho)$

or (iv) $\lambda = \lambda(p, \rho)$

[Prove it. Need chain rule for (iii) and (iv).]

Assuming one of these latter conditions is true, we get:

$$\boxed{\frac{D}{Dt} \left(\frac{\bar{\omega} \cdot \nabla \lambda}{\rho} \right) = 0} \quad \underline{\text{Ertel's Potential Vorticity Theorem}}$$

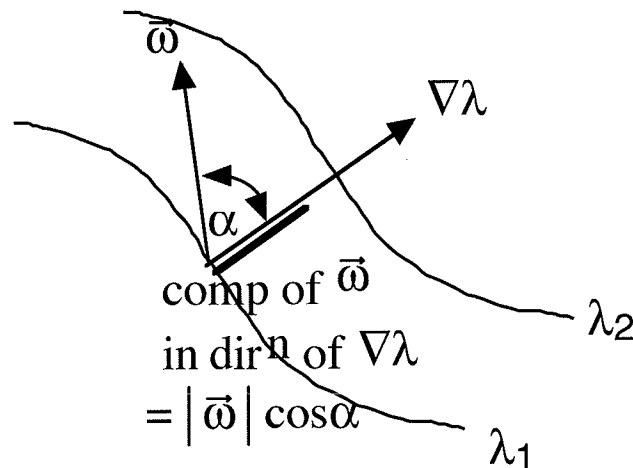
or,

$$\boxed{\frac{\bar{\omega} \cdot \nabla \lambda}{\rho} = \text{const}} \quad \text{for a parcel (dif. consts for dif parcels).}$$

$\frac{\bar{\omega} \cdot \nabla \lambda}{\rho}$ is Ertel potential vorticity

$$= \frac{|\nabla\lambda|}{\rho} \left[|\vec{\omega}| \cos \alpha \right] \quad \text{where } \alpha \text{ is angle btw } \vec{\omega} \text{ and } \nabla\lambda .$$

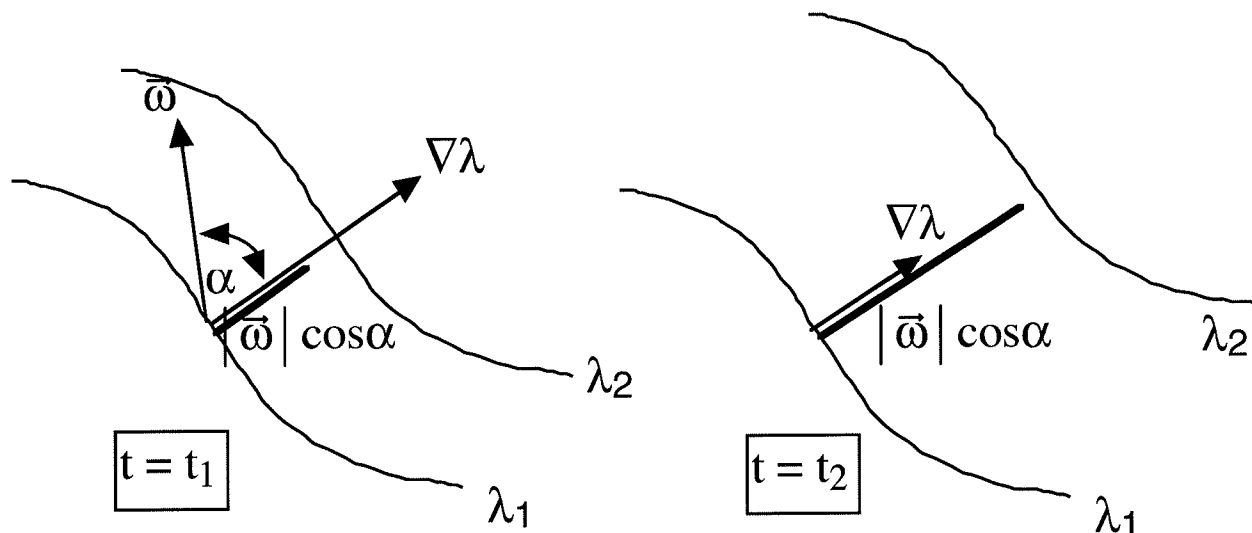
Product of magnitude of gradient of conserved scalar and comp of vorticity in direction of that gradient ($\div \rho$) is const for a parcel.



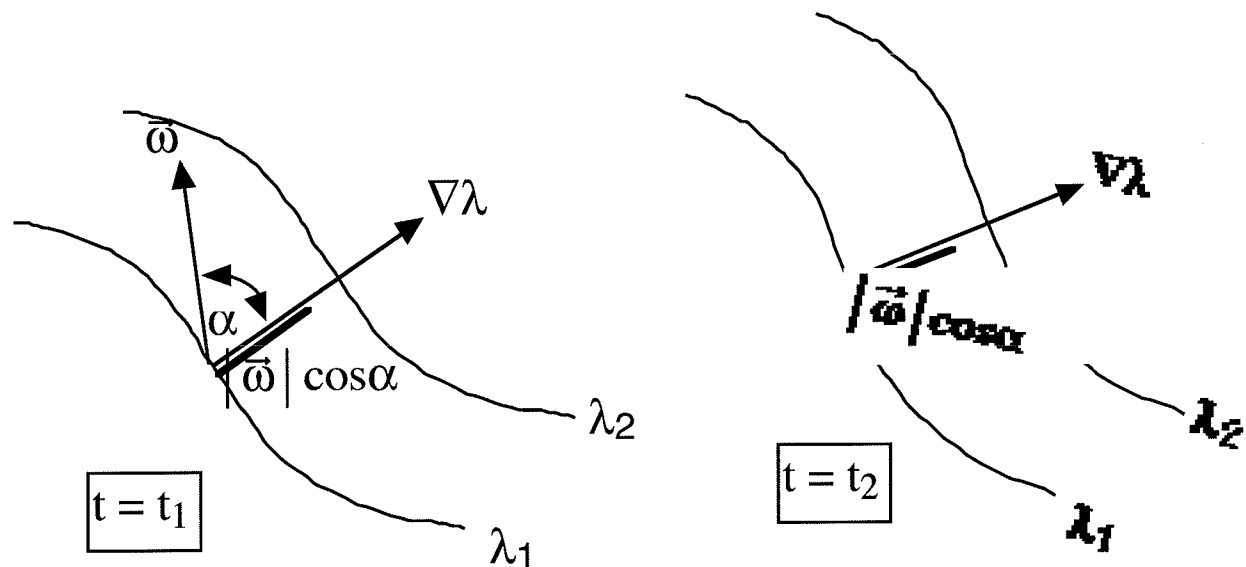
[Comp of $\vec{\omega} \perp \nabla\lambda$ is irrelevant to thm.]

e.g. to isolate role of ρ , consider case where parcel moves downward \therefore parcel compresses, so $\rho \uparrow$ so Ertel's thm says vorticity \uparrow (provided we ignore changes in $|\nabla\lambda|$). What happened? Mass was brought in closer to rot axis so vort \uparrow .

e.g., suppose flow pulls λ sfcs apart. In that case $|\nabla\lambda| \downarrow$
 \therefore comp of $\vec{\omega}$ in dirⁿ of $\nabla\lambda \uparrow$. Increase is due to stretching:



e.g., suppose flow rotates λ sfc's w/out changing magnitude of $\nabla\lambda$:



- so flow rotates $\nabla\lambda$ vector w/out changing magnitude of $\nabla\lambda$ so magnitude of comp of $\vec{\omega}$ in dirⁿ of $\nabla\lambda$ stays the same. [in this example vort vector would also rotate w/ flow]